

Experimental validation of a vanadium redox flow battery model for state of charge and state of health estimation

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Abstract This study presents a vanadium redox flow battery model that considers the most important variables that have a crucial role in the performance of the system. A complete model divided in an electrochemical, thermal, hydraulic and voltage submodels is presented. The analytic analysis of the model is carried out to reduce the system order according to some conservation laws. Based on this analysis, a subsequent calibration of the model parameters is developed using real experimental data. The validation is performed comparing the real measured voltage and the one estimated with the model. To calibrate the model an algorithm based on the implementation of a particle swarm optimizer is used. Results obtained in both short and long-term operation are presented, in order to compare and validate if the model can be used for both state of charge and state of health estimation.

Keywords: Redox flow battery, Nonlinear model, State of Charge, State of Health, Particle Swarm Optimization

1 INTRODUCTION

Among all types of energy storage systems (ESS), vanadium redox flow batteries (VRFB) stand out for their high efficiency (70-90%), long life cycle, safety and the possibility to decouple both energy and power sizes [4].

One of the main challenges in the study and analysis of a VRFB consists on the correct characterization of its behaviour through mathematical models. The importance of developing mathematical models for VRFB is reflected in the literature [5] [17], where it is possible to find one of the most accepted electrochemical models developed by Skyllas-Kazacos, who pioneered the use of vanadium in redox flow batteries [19].

Models can be classified as static or dynamic, depending on whether the effect of time is taken into account, and as distributed or lumped parameter models, depending on the space dimension. Tools such as COMSOL are used to analyze and develop distributed models. However, for control purposes, distributed models are not used due to its complexity, relevant computational cost, resources and time [3]. In counterpart, vast majority of works concerning VRFB control problems, use lumped parameters models.

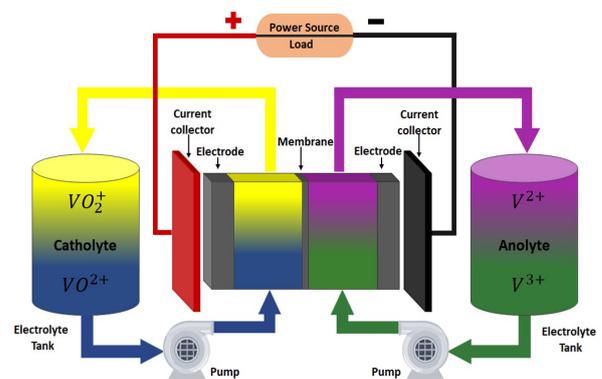


Figure 1: Vanadium redox flow battery scheme with its main components.

Most of these lumped parameter models use different hypothesis and assumptions that simplify their complexity. Although a vast majority of these assumptions may not be entirely realistic, they help in the understanding, formulation and analysis of the model. This is the case of the mentioned Skyllas-Kazacos electrochemical models, where the same flow rate is considered in both parts of the system, same species concen-

32 tration inside the cell and tanks, as well as a constant
33 temperature to estimate the open circuit voltage (OCV)
34 by the Nernst equation.

35 In general, most of the works found in the litera-
36 ture related to VRFB modeling, study and analyze the
37 behaviour of different variables as the evolution of the
38 species concentration, voltage, temperature, pressure,
39 state of charge (SOC) or state of health (SOH). The
40 evolution of species concentrations can be found in lot
41 of studies, from lesser to greater complexity. Thus,
42 Skyllas-Kazacos presents in [18] a model that expresses
43 the dynamics of the vanadium species considering the
44 different ion transport mechanisms such as diffusion,
45 migration and convection. This model has been im-
46 proved in some works introducing the effect of other
47 phenomena, such as the oxidation of vanadium V^{2+} due
48 to the presence of air inside the negolyte tank [11].
49 These models only consider the participation of vana-
50 dium species. However, there are other species that can
51 be considered such as the water and hydrogen ions.
52 The inclusion of these species not only improves the
53 better understanding of the electrochemical part, but
54 also allows to consider other important effects, such as
55 the Donnan potential in the computation of the voltage
56 [10]. The dynamics of the water crossover phenomena
57 are analyzed in [15], while those of the hydrogen ions
58 have also been studied [16].

59 Regarding voltage computation, most of the works
60 are focused on the determination and analysis of the
61 different overpotentials and phenomena that affect its
62 computation. The OCV is discussed in [7], showing the
63 discrepancy between the common models and experi-
64 mental data, presenting a complete form of the Nernst
65 equation considering the proton activity and Donnan
66 potential. Other works such as [9] and [13] are focused
67 on the computation of the concentration and activation
68 overpotentials.

69 With respect to the thermal model, different works
70 presented in the literature study the temperature dy-
71 namics in some parts of the system, those inside the
72 cell being specially interesting. Thus, a thermal model
73 is shown in [24], that allows to compute the tempera-
74 ture inside the cell taking into account the temperature
75 in the pipes and tanks. The inclusion of the dynam-
76 ics due to self-discharge reactions is presented in [20]
77 while the analysis of the shunt current effect is investi-
78 gated in [23].

79 Pressure is another important variable to take into
80 account that is directly related to the flow rates of the
81 system, which have an important role in control pur-
82 poses such as optimal control. Therefore, some stud-
83 ies analyze the behaviour of the pressure inside a VRFB
84 system as can be observed in [22].

85 Among all variables of interest regarding the per-
86 formance of a VRFB, state of charge (SOC) and state of

87 health (SOH) stand out as the most important. There
88 exist different definitions for these variables, according
89 to the parameters and variables that are considered for
90 their computation [14]. However, the vast majority of
91 these definitions present the SOC and SOH as a func-
92 tion of the species concentrations, which makes it nec-
93 essary to choose an appropriate electrochemical model
94 [6].

95 Usually, most of the studies that present an elec-
96 trochemical model for SOC and SOH computation are
97 only validated theoretically or by means of simulations.
98 However, experimental validation is usually limited to
99 the estimation of one of these variables, and under
100 some assumptions that can differ from a real battery
101 behaviour. Moreover, for this cases, it is more difficult
102 to find studies that present a general model that com-
103 prises the most important variables described in a sin-
104 gle model.

105 Thus, in this work, a complete model that considers
106 the most important effects of a VRFB is presented. The
107 model can be divided into four submodels, according to
108 the most important variables, which are: electrochemi-
109 cal, voltage, thermal and hydraulic ones. The main ob-
110 jective of this work is to understand the dynamics of
111 each one of these subsystems in an efficient way, pre-
112 senting the main equations that describe the different
113 effects, and relate them with the other submodels. With
114 these dynamics, the most important performance indi-
115 cators such as the SOC, SOH, capacity and efficiencies
116 are defined. A general model that can deal with short
117 and long term scenarios is presented in order to see if
118 it is suitable for SOC and SOH. In order to validate the
119 model, an experimental setup has been developed to es-
120 timate these two parameters calibrating the model with
121 voltage and current measurements obtained from dif-
122 ferent experiments in short and long term campaigns.
123 In order to calibrate the model, there exist different al-
124 gorithms and methods that can be used, being the par-
125 ticle swarm optimization (PSO) a good choice to esti-
126 mate a certain number of parameters whose value is
127 only known within a particular range [26]. Using the
128 PSO technique it is possible to calibrate the model us-
129 ing experimental data extracted from a VRFB that has
130 been designed, consisting on a cell with two little tanks,
131 presenting a similar aspect to the VRFB scheme shown
132 in Figure 1.

133 The paper is organized as follows. Section 2
134 presents the formulated model, while in Section 3, a
135 conservation study is developed in order to see if the
136 mass and charge conservation principles are fulfilled for
137 the model dynamics presented. Section 4 presents the
138 experimental validation of the model, describing the
139 experimental setup and the different results obtained
140 from the model calibrated using the PSO algorithm. Fi-
141 nally, Section 5 presents the main conclusions.

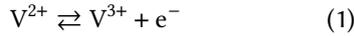
142 2 MODEL FORMULATION

143 The model can be formulated in different submodels ac-
144 cording to the properties and characteristics of the state
145 variables. The four submodels into which the general
146 model has been divided are the electrochemical, the
147 thermal, the hydraulic and the voltage ones. All sub-
148 models interact with the others, sharing some parame-
149 ters and variables to define the most important perfor-
150 mance indicators, typical of any energy storage system.

151 2.1 Electrochemical model

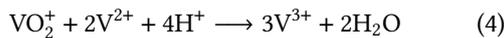
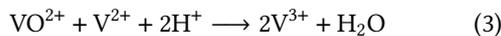
152 Among the different electrochemical models that ex-
153 ist for VRFB, the dynamic model proposed by Skyllas-
154 Kazacos is the most widely used and accepted [18].
155 Based on this model, it is possible to express the be-
156 haviour of the vanadium species inside the system, con-
157 sidering the most important phenomena such as the re-
158 dox reactions and the ion transport due to ion crossover
159 methods.

160 On the one hand, the redox reactions that take place
161 inside the cell/stack of a VRFB are:

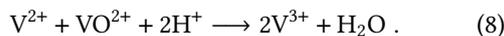
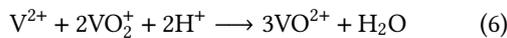


162 where V^{2+} and V^{3+} are the vanadium species in the
163 negative side (1), and VO_2^+ and VO^{2+} are the vanadium
164 species in the positive side (2), which are commonly ex-
165 pressed by the oxidation state as V^{4+} and V^{5+} , respec-
166 tively.

167 On the other hand, the ion crossover methods ap-
168 pear due to the self-discharge reactions that occur on
169 both sides of the system. The mass balance due to these
170 reactions in both positive and negative electrodes can
171 be described separately. On the one hand, in the nega-
172 tive electrode the self-discharge reactions occurs due to
173 the presence of VO_2^+ and VO_2^+ that have moved across
174 the membrane. These species are reduced by V^{2+} to be-
175 come V^{3+} :



176 On the other hand, the same behaviour occurs in
177 the positive electrode, where vanadium species V^{2+} and
178 V^{3+} cross the membrane and are oxidized by VO_2^+ to
179 VO^{2+} :



180 Skyllas-Kazacos considers several assumptions that
181 are the following [1]:

- 182 • Electrolytes are perfectly mixed.
- 183 • The total number of moles of the species remains
184 constant.
- 185 • Vanadium concentrations are uniformly dis-
186 tributed in the cell and tanks, without consider-
187 ing the hydrogen ions and water.
- 188 • Electrolyte flow rate is the same in both half-cells.
- 189 • Electrolyte temperature remains constant.
- 190 • Self-discharge reactions occurs instantaneously
191 by ion crossover.
- 192 • Vanadium and hydrogen ions are transported
193 through the membrane by diffusion. Other trans-
194 port mechanisms of water and ions such as mi-
195 gration, convection or osmosis are neglected.
- 196 • Diffusion coefficients are constant.
- 197 • Gas evolution or other side reactions during
198 charging/discharging neglected.

199 Considering the above assumptions, the general ex-
200 pression that describes the concentration evolution of
201 the species is:

$$\varepsilon v^e \frac{\partial c_i^c}{\partial t} = q \cdot (c_i^t - c_i^c) \pm s^e \frac{j}{F} + d_i, \quad (9)$$

202 where ε is the electrode porosity, v^e is the electrode vol-
203 ume and c_i^c is the concentration of vanadium species i
204 inside the cell, expressed in $\text{mol}\cdot\text{m}^{-3}$.

205 The electrode volume v^e depends only on the di-
206 mensions of the electrode and the number of equal cells
207 of the stack:

$$v^e = N \cdot h^e \cdot w^e \cdot l^e, \quad (10)$$

208 N being the number of cells, and h^e , w^e and l^e the
209 height, thickness and length, respectively, in m . These
210 dimensions have been defined in this work considering
211 the orientation shown in Figure 2.

212 As can be noticed, the cell concentration expression
213 (9) can be splitted in three parts. The first term defines
214 the change of species due to the inlet and outlet flux of
215 electrolyte, generated by the electrolyte flow rate q in
216 $\text{m}^3 \cdot \text{s}^{-1}$. The second term implies the role of the redox
217 reactions (1) - (2) due to the presence of the electrons
218 generated for the existing current between the collec-
219 tors. This effect represents the current flux perpendic-
220 ular to the cell, and is described by a current density
221 j , expressed in $\text{A}\cdot\text{m}^{-2}$, F that is the Faraday constant
222 ($96845.33 \text{ C}\cdot\text{mol}^{-1}$) and the electrode surface s^e which
223 depends on the electrode height and length:

$$s^e = h^e \cdot l^e. \quad (11)$$

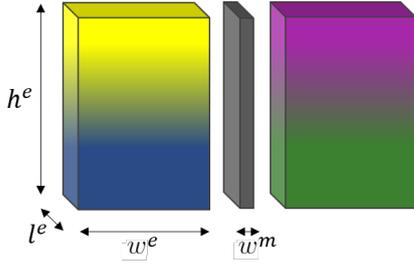


Figure 2: 3D Scheme of a single cell composed by two electrodes and a membrane. h^e : electrode and membrane height, l^e : electrode and membrane length, w^e : electrode thickness, w^m : membrane thickness.

Finally, last term d_i is the total flux of species i across the membrane, which is the ion crossover effect and, in this work, it is only associated with the diffusion. The diffusion phenomenon occurs because the membranes allow the crossover of vanadium species from one side of the cell to the other. This movement from high to low concentration areas is driven by concentration gradients, and therefore depends on each vanadium species, modelled to obey Fick's law:

$$d_i = \frac{s^m}{w^m} \cdot (\alpha_{2_i} k_2 c_2^c + \alpha_{3_i} k_3 c_3^c + \alpha_{4_i} k_4 c_4^c + \alpha_{5_i} k_5 c_5^c), \quad (12)$$

where s^m is the membrane surface in contact with the electrolyte, which according to Figure 2 can be computed as s^e , α_{k_i} is the mass balance between the particular vanadium species i with the other species k , and k_i is the diffusion coefficient of each vanadium species, expressed in $\text{m}^2 \cdot \text{s}^{-1}$.

In most of the studies, the diffusion coefficients k_i are assumed to be constant and invariant. However in this work a variable obeying the Arrhenius equation have been considered [21]:

$$k_i = A_i e^{\frac{-E_a}{RT^c}}, \quad (13)$$

A_i being the pre-factor, which is a constant that depends on the membrane material and is different for each vanadium specie, E_a the activation energy of the vanadium species ($17340 \text{ J} \cdot \text{mol}^{-1}$), which is assumed to be the same for all species [27], R is the universal gas constant ($8.314 \text{ J} \cdot \text{mol}^{-1} \cdot \text{K}^{-1}$) and T^c is the electrolyte temperature inside the cell/stack. Therefore, each species has a different pre-factor that depends on the membrane composition.

For the tank dynamics, as they are isolated parts of the system whose only function is to store the electrolyte, the only important term is the inlet and outlet flux of electrolyte:

$$v^t \frac{\partial c_i^t}{\partial t} = q \cdot (c_i^c - c_i^t), \quad (14)$$

where v^t is the electrolyte volume inside the tank and c_i^t is the concentration of species in the tank, sharing the same units of $\text{mol} \cdot \text{m}^{-3}$ as the cell concentration. As there is not any flux of current going through the tanks, and species cannot escape from them, the second and third terms of (9) are neglected.

Considering the cell and tank dynamic expressions formulated by (9) and (14), some modifications have been made respect to the model presented by Skyllas-Kazacos [18] in order to make it more realistic. They are:

- Electrolyte flow rate can be different in both half-cells. Thus, a clear distinction between the electrolyte circuits is made, distinguishing the negative (-) from the positive (+) side.
- Electrolyte temperature varies and it is computed by a thermal model closely related to the electrochemical one.
- Hydrogen ions and water are considered. They are usually neglected, but have an important role in the system operation.

As this work pretends to be a tool for the understanding and developing of a possible model, the vanadium, water and hydrogen ion models are presented separately. All of them are formulated in the state space domain to facilitate the understanding and comparison.

2.1.1 Vanadium electrochemical model

Considering expressions (9) and (14), the vanadium electrochemical model can be formulated in the state space notation as:

$$\dot{\mathbf{x}} = \mathbf{A}\mathbf{x} + \mathbf{B}_- \mathbf{x} \cdot q_- + \mathbf{B}_+ \mathbf{x} \cdot q_+ + \mathbf{b} \cdot j, \quad (15)$$

where q_- and q_+ are the flow rates of negolyte and posolyte, j is the current density and the state vector, \mathbf{x} , is defined as $\mathbf{x} = [c_2^c \ c_3^c \ c_4^c \ c_5^c \ c_2^t \ c_3^t \ c_4^t \ c_5^t]^T$; where c_i^p stands for the concentration of vanadium species i in space p , with $p = \{c, t\}$ meaning cell and tank, respectively. Matrix \mathbf{A} is directly related to the diffusion d_i , matrices \mathbf{B}_- and \mathbf{B}_+ with the flow rates effect and vector \mathbf{b} is related to the presence of electrons due to the redox reaction. All matrices and vectors appear summarized in the Appendix section.

Using this formulation, j , q_- and q_+ are the input variables which are assumed to be known, being easily measured and controlled.

2.1.2 Hydrogen electrochemical model

The same procedure can be developed for the hydrogen ions model. In that case, it can be represented as:

$$\dot{\mathbf{x}}_H = \mathbf{A}_H \mathbf{x}_H + \mathbf{B}_H \mathbf{x}_H \cdot q_- + \mathbf{B}_H \mathbf{x}_H \cdot q_+ + \mathbf{B}_d \mathbf{x}_H + \mathbf{b}_H \cdot j, \quad (16)$$

where \mathbf{x}_H is the new state vector of proton concentrations considering the cell and tanks and is defined as $\mathbf{x}_H = [c_{H_+}^c \ c_{H_+}^t \ c_{H_+}^c \ c_{H_+}^t]^\top$. Matrix \mathbf{A}_H expresses the diffusion due to the concentration difference of both sides. Matrices \mathbf{B}_{H_+} and \mathbf{B}_H are related to the inlet and outlet fluxes from cell and tanks and matrix \mathbf{B}_d is related to the proton diffusion due to the self-discharge reactions, which is related to the vanadium species x . Finally, vector \mathbf{b}_H expresses the variation of protons due to the charging/discharging current density j where, according to the redox reactions (1) - (2), it is possible to see how in the negative part, the protons do not appear, while in the positive one, the redox reaction involves 2 hydrogen ions H^+ for each electron e^- .

2.1.3 Water electrochemical model

Another important species to consider is the water. In this case, the electrochemical expressions that describe its evolution inside the system can be formulated in state space as:

$$\dot{\mathbf{x}}_{H_2O} = \mathbf{A}_{H_2O} \mathbf{x} + \mathbf{B}_{H_2O} \cdot \mathbf{x}_{H_2O} \cdot q_- + \mathbf{B}_{H_2O} \cdot \mathbf{x}_{H_2O} \cdot q_+ + \mathbf{b}_{H_2O} \cdot j, \quad (17)$$

where \mathbf{x}_{H_2O} is the state vector of water concentrations considering the cell and tanks, and it is defined as $\mathbf{x}_{H_2O} = [c_{H_2O}^c \ c_{H_2O}^t \ c_{H_2O}^c \ c_{H_2O}^t]^\top$. The matrix \mathbf{A}_{H_2O} contains the diffusion coefficients of the vanadium species and is related to the self-discharge diffusion. Matrices \mathbf{B}_{H_2O} and \mathbf{B}_{H_2O} are related to the inlet and outlet fluxes and are exactly equal to the protons matrices, and vector \mathbf{b}_{H_2O} is related to the variation of water in the positive part due to the current density.

2.2 Performance parameters

Among the many parameters that can be established depending on the variable or characteristic to be analyzed, it is worth highlighting some of them that are commonly studied and used for any type of ESS. These main indicators are the SOC and SOH.

2.2.1 State of charge

The state of charge is one of the most common indicators of a battery, that is used to denote the capacity that is currently available. It is directly related to the amount of species inside the global system. In particular, the SOC of a VRFB is computed as the ratio of active species V^{2+} and VO_2^+ with respect to the total concentration of species in each side. Thus, it should be noted that there exist two different SOC, defined as (SOC₋) and (SOC₊) to represent the SOC in the negolyte and posolyte, respectively.

As the volume of the tanks is generally much greater than that of the cell/stack, the role of the cell concen-

tration is neglected to calculate the SOC in most studies. Moreover, many of them assume that the concentrations on both sides of the system are the same and evolve equally with time, concluding that there exist only one definition of SOC, being exactly the same for both parts:

$$SOC = \frac{c_2}{c_2 + c_3} = \frac{c_5}{c_4 + c_5}. \quad (18)$$

Although these assumptions allow to simplify the VRFB model, they have not been considered since the purpose of this study is to develop a complete and realistic model that includes the main operating phenomena, as well as degradation. Therefore, considering the effect of the species transport mechanism, it is essential to differentiate between negolyte and posolyte concentrations, obtaining the following expressions:

$$SOC_- = \frac{c_2}{c_2 + c_3} \quad (19)$$

363

$$SOC_+ = \frac{c_5}{c_4 + c_5}. \quad (20)$$

Furthermore, considering the role of the cell concentration, that has influence on the system voltage and may be different from that of the tanks, expressions (19) and (20) must be reformulated as:

$$SOC_- = \frac{c_2^t v_-^t + c_2^c v^c}{c_2^t v_-^t + c_2^c v^c + c_3^t v_-^t + c_3^c v^c} \quad (21)$$

368

$$SOC_+ = \frac{c_5^t v_+^t + c_5^c v^c}{c_4^t v_+^t + c_4^c v^c + c_5^t v_+^t + c_5^c v^c}. \quad (22)$$

As can be noticed from (21) and (22), the SOC has been recomputed as a ratio of mass species, taking into account the volumes of the cell and tanks. This expression is more realistic and considers all vanadium species concentrations along the system.

However, it is important to note that the total amount of energy accumulated in the system will be given by the electrolyte with less SOC, in order to restrict the maximum energy that can be extracted. Accordingly, the correct expression of the SOC must be defined as:

$$SOC = \min(SOC_-, SOC_+). \quad (23)$$

2.2.2 State of health

The state of health can be computed as the ratio between the amount of mass available in a certain moment with respect to the original one. In this case, the SOH has not been extensively analyzed in the literature, since, as has been commented, many of the studies do not consider the mass transport mechanisms. Nevertheless, it is one of the most important parameters since it is relevant for the life cycle of a battery.

389 For the VRFB, following with the definition pre-
 390 sented for the SOC in (23), the SOH can be computed
 391 as the minimum amount of moles available between the
 392 two parts of the system, with respect to the ideal one:

$$\text{SOH} = \frac{\min(m_-, m_+)}{m^*}, \quad (24)$$

393 where m_- and m_+ define, respectively, the total number
 394 of moles, in the negolyte and posolyte parts, which can
 395 be computed by means of the following expressions:

$$m_- = v^c \cdot c_-^c + v_-^t \cdot c_-^t \quad (25)$$

396

$$m_+ = v^c \cdot c_+^c + v_+^t \cdot c_+^t, \quad (26)$$

397 c_- being the sum of V^{2+} and V^{3+} species concentrations
 398 corresponding to the negolyte and c_+ the sum of VO^{2+}
 399 and VO_2^+ for the posolyte part.

400 The other variable of (24) is the ideal mole distribu-
 401 tion m^* , which can be defined from the total number of
 402 vanadium moles of the system as:

$$m^* = \frac{m}{2}, \quad (27)$$

403 where m denotes the vanadium moles of the system,
 404 which can be computed as the sum between both parts:

$$m = m_- + m_+ \quad (28)$$

405 2.3 Thermal model

406 Following the principle of conservation of energy, a dy-
 407 namic thermal model can be formulated based on en-
 408 ergy balance equations for the different parts of the
 409 system. Thus, temperature inside both tanks, cell and
 410 pipes can be considered. Regarding the heat generated
 411 from the pumps, it is convenient to differentiate be-
 412 tween the temperature inside the pipes in two different
 413 parts, as depicted in Figure 3.

414 Among different thermal dynamic models for VRFB
 415 based on conservation of energy that appear in the lit-
 416 erature, Tang et al. presents a model that considers the
 417 temperatures displayed in Figure 3, and takes into ac-
 418 count the effect of pump heat and the self-discharge re-
 419 actions due to the ion crossover mechanisms [24]. How-
 420 ever, despite being a fairly realistic model, different as-
 421 sumptions have been taken into account:

- 422 • Temperature is uniformly distributed in each part
 423 of the system.
- 424 • Heat is transferred in each part between the elec-
 425 trolyte and the surrounding air.
- 426 • Cell heat is generated by the cell resistance and
 427 electrochemical and self-discharge reactions.

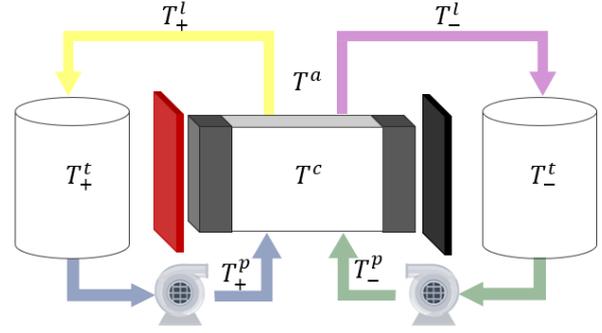


Figure 3: Temperature measurements in the different parts of the system. T_+^t : posolyte tank, T_-^t : posolyte line, T_+^p : posolyte pump line, T_-^p : negolyte tank, T_-^t : negolyte line, T_-^p : negolyte pump line, T^c : cell and T^a : air temperature.

428 One of the novelties of this work is the introduction
 429 of the heat generated due to the chemical reactions,
 430 considering their entropy. This new consideration adds
 431 more reality to some of the accepted thermal models
 432 presented in different works [25].

433 Similarly to the electrochemical model, the temper-
 434 ature inside each part T_i , changes according to the inlet
 435 and outlet flux of electrolyte with the contiguous part
 436 T_o . This change of temperature can be formulated as:

$$v c_p \rho \frac{\partial T_i}{\partial t} = q_k c_p \rho \cdot (T_o - T_i), \quad (29)$$

437 where q_k is the electrolyte flow rate in the posolyte and
 438 negolyte sides, expressed in $\text{m}^3 \cdot \text{s}^{-1}$, c_p is the specific
 439 heat capacity of the electrolyte in $\text{J} \cdot \text{kg}^{-1} \cdot \text{K}^{-1}$ and ρ is the
 440 electrolyte density expressed in $\text{kg} \cdot \text{m}^{-3}$.

441 Considering (29) and the previous assumptions, it
 442 is possible to express the temperature inside the cell T^c
 443 as:

$$\begin{aligned} 2v^c c_p \rho \frac{\partial T^c}{\partial t} &= q_+ c_p \rho \cdot (T_+^p - T^c) + q_- c_p \rho \cdot (T_-^p - T^c) \\ &+ u^c s^c \cdot (T^a - T^c) + (s^e j)^2 r \\ &+ s^e \frac{j}{F} \cdot (\Delta S_+ - \Delta S_-) T^c + d_T, \end{aligned} \quad (30)$$

444 where T_+^p and T_-^p are the temperatures at the pump out-
 445 let in the positive and negative sides, respectively in K,
 446 u^c is the overall heat transfer coefficient of the cell in
 447 $\text{W} \cdot \text{m}^{-2} \cdot \text{K}^{-1}$, s^c is the cell surface in m^2 , T^a is the sur-
 448 rounding air temperature and r is the cell resistance in
 449 Ω . The first and second terms represent the change of
 450 temperature according to (29) for the positive and neg-
 451 ative parts, respectively. The third term is related to the
 452 heat transferred between the cell surface and the sur-
 453 rounding air. The fourth term defines the heat gener-
 454 ated inside the cell in form of Ohmic losses due to the

455 Joule effect. Fifth term represents the change of tem-
 456 perature due to entropy. Finally, d_T defines the changes
 457 of temperature due to the self-discharge reactions. This
 458 last term is related to the changes of enthalpy ΔH of the
 459 different reactions and can be modeled as:

$$d_T = \frac{s^m}{w^m} [k_2 c_2^c \cdot (-\Delta H_{(3)}) + k_3 c_3^c \cdot (-\Delta H_{(4)}) + k_4 c_4^c \cdot (-\Delta H_{(6)}) + k_5 c_5^c \cdot (-\Delta H_{(7)})], \quad (31)$$

460 where $\Delta H_{(i)}$ is the change of enthalpy for the self-
 461 discharge reactions. The enthalpies of the different
 462 species are constant and can be found in the litera-
 463 ture. Considering their values it is possible to obtain the
 464 change of enthalpy of the self-discharge reactions (3 -
 465 4) and (6 - 7). Table 1 presents the value of the species
 466 and self-discharge reactions enthalpies.

467 Once the cell temperature evolution has been clar-
 468 ified, it is possible to develop the same analysis for the
 469 other parts of the system. It is important to notice that
 470 the expressions for the negative and positive part of the
 471 system are equal, due to the fact that can be seen as
 472 symmetrical with respect to the cell. In the vast major-
 473 ity of studies, they are considered equal, because they
 474 can be seen as symmetrical, assuming same dimensions
 475 and flow rates in both parts. However, as one of the
 476 main objectives of this research is to reach a general
 477 model that can be used in a wide range of operational
 478 conditions, they have been considered separately.

Table 1: Enthalpies of the species at 298.15 K and change of enthalpies of self-discharge reactions.

Species / Reactions	ΔH (kJ·mol ⁻¹)
V ²⁺	-226.0
V ³⁺	-259.0
VO ²⁺	-486.8
VO ₂ ⁺	-649.8
H ₂ O	-285.8
H ⁺	0
(3)	-91.2
(4)	-246.8
(6)	-220.0
(7)	-64.0

479 On the one hand, the expression for the temper-
 480 ature of the line that connects the cell with the tank
 481 without the pump, T_k^l , is:

$$v_k^l c_p \rho \frac{dT_k^l}{dt} = q_k c_p \rho \cdot (T^c - T_k^l) + u^l s_k^l \cdot (T^a - T_k^l), \quad (32)$$

482 where v_k^l is the volume of this part of the pipe, u^l is the
 483 overall heat transfer coefficient of the pipe and s_k^l is the
 484 pipe surface of that section.

485 For the case of the tank temperature, its expression
 486 is:

$$v_k^t c_p \rho \frac{dT_k^t}{dt} = q_k c_p \rho \cdot (T_k^l - T_k^t) + u^t s_k^t \cdot (T^a - T_k^t), \quad (33)$$

487 T_k^t being the temperature inside the posolyte or ne-
 488 golyte tank, u^t its heat transfer coefficient and s_k^t its
 489 surface. Finally, the temperature in the pipe line where
 490 the pump is located, T_k^p , can be expressed as:

$$v_k^p c_p \rho \frac{dT_k^p}{dt} = q_k c_p \rho \cdot (T_k^t - T_k^p) + u^p s_k^p \cdot (T^a - T_k^p) + W_k^p, \quad (34)$$

491 where v_k^p is the pipe volume of that section, u^p is the
 492 heat transfer coefficient, which is equal to u^l as both are
 493 pipes, s_k^p is its surface and W_k^p is the heat generated by
 494 the pump, which is formulated in the hydraulic model
 495 section.

496 Finally, the thermal model can be expressed in the
 497 state space formulation using same notation that the
 498 one used for the electrochemical model to summarize
 499 all previous expressions:

$$\dot{T} = \mathbf{A}_T T + \mathbf{B}_{T,+} T \cdot q_+ + \mathbf{B}_{T,-} T \cdot q_- + \mathbf{B}_H x + \mathbf{b}_T \cdot j^2 + \mathbf{c}_T \cdot T^a + \mathbf{w}^p, \quad (35)$$

500 where T is the temperature state vector defined as $T =$
 501 $[T^c \ T_+^l \ T_+^t \ T_+^p \ T_-^l \ T_-^t \ T_-^p]^\top$, matrices \mathbf{A}_T , $\mathbf{B}_{T,+}$
 502 and $\mathbf{B}_{T,-}$ are related to the inlet outlet temperature and
 503 air surrounding heat, matrix \mathbf{B}_H with the change of en-
 504 thalpies, vector \mathbf{b}_T is related to the ohmic losses, vector
 505 \mathbf{c}_T with the room temperature and \mathbf{w}^p with the heat
 506 generated on the pumps.

507 2.4 Hydraulic model

508 Based on different works found in the literature that
 509 describe in detail the hydraulic part of a VRFB system,
 510 a simple model is presented that only considers the ele-
 511 ments that present a high contribution in terms of pres-
 512 sure drop.

513 This hydraulic model, like the thermal one, uses the
 514 principle of conservation of energy. In this way, the
 515 power generated by the pumps, W^p , can be computed
 516 as the product between the pressure drop inside the
 517 system, Δp , and the electrolyte flow rate q :

$$W^p = \Delta p \cdot q. \quad (36)$$

518 It should be noted that some other heat sources,
 519 such as the friction of the moving elements of the pump
 520 with the electrolytes, are neglected.

521 There are three different elements that have been
 522 chosen as main contributors of pressure drop, which are
 523 the pipes, the flow frames and the electrodes. The de-
 524 termination of the pressure drop in each part is based
 525 on the laws of fluid mechanics. The same procedure to

526 the one developed in the thermal model to differenti-
527 ate between positive and negative parts has been used,
528 using a variable k to distinguish between both parts.

529 On the one hand, the pressure drop due to the elec-
530 trolyte circulating along the hydraulic pipes, Δp_k^l , can
531 be computed as:

$$\Delta p_k^l = f \frac{l_k^l}{d_k^l} \frac{\rho v_k^2}{2} \quad (37)$$

532 where f is the friction loss factor, l_k^l is the length of the
533 hydraulic pipe line in m, d_k^l is the hydraulic diameter of
534 the pipes in m, and v_k is the velocity of the electrolyte,
535 which is expressed in $\text{m} \cdot \text{s}^{-1}$ and can be computed as:

$$v_k = \frac{q_k}{s_k^l} \quad (38)$$

536 q_k being the electrolyte flow rate and s_k^l the cross-
537 section of the pipe in m^2 .

538 The pressure drop in the flow frames follows the
539 same law and can be defined with the same expression:

$$\Delta p_k^f = f \frac{l^f}{d^f} \frac{\rho v_k^2}{2} \quad (39)$$

540 where l^f is the length of the electrolyte flow frames and
541 d^f is its hydraulic diameter, which is assumed to be the
542 same for negolyte and posolyte parts, as the cells usu-
543 ally have same dimensions and materials.

544 Finally, the pressure drop in the electrodes is calcu-
545 lated using Darcy's law in porous media by:

$$\Delta p_k^e = \frac{\mu q_k K_{ck} (1 - \varepsilon)^2}{d_f^2 \varepsilon^3} \frac{h^e w^e}{l^e} \quad (40)$$

546 where μ is the viscosity of the fluid in $\text{Pa} \cdot \text{s}$, K_{ck} is the
547 Kozeny-Carman constant and d_f is the mean diameter
548 of the electrode fibers expressed in m.

549 Then, the total power supplied by the pumps to
550 overcome the pressure drop in the system can be com-
551 puted as:

$$W_k^p = \left(\frac{l_k^l}{d_k^l} \frac{\rho v_k^2}{2} + f \frac{l^f}{d^f} \frac{\rho v_k^2}{2} + \frac{\mu q_k K_{ck} (1 - \varepsilon)^2}{d_f^2 \varepsilon^3} \frac{h^e w^e}{l^e} \right) q_k \quad (41)$$

552 2.5 Voltage model

553 Voltage computation is one of the most important
554 parts, since it allows to determine the power and energy
555 stored in the system. Moreover, it is an analog variable
556 that in practice can be measured by means of analog or
557 digital instruments. In this way, it stands out as one of
558 the most used variables in the validation of any VRFB
559 model.

560 Ideally, the voltage of any electrochemical system
561 can be split in two parts, which are the ideal voltage,
562 and some losses due to different effects called overpo-
563 tentials. On the one hand, the ideal voltage only de-
564 pends on the species and the chemical reaction that
565 takes place. This ideal voltage produced only by the
566 effect of the chemical reaction is called open circuit
567 voltage (OCV) and can be computed by means of the
568 Nernst equation. On the other hand, the charging or
569 discharging current across the RFB system implies the
570 appearance of the different overpotentials such as the
571 activation, ohmic and concentration ones. Therefore,
572 the total voltage of a RFB, parameterized as E , can be
573 expressed as the sum of the OCV and the overpoten-
574 tials:

$$E = E^{OCV} + \eta^a + \eta^o + \eta^c \quad (42)$$

575 where E^{OCV} is used in this work to represent the OCV,
576 and η^a , η^o and η^c to define the activation, ohmic and
577 concentration overpotentials, respectively.

578 2.5.1 OCV formulation

579 The computation of the OCV for any general electro-
580 chemical system is based on the Nernst equation.

581 For the case of a VRFB the well-known general ex-
582 pression to compute the OCV is:

$$E^{OCV} = E^\theta + \frac{RT^c}{F} \ln \left(\frac{c_2 \cdot c_5 \cdot c_{H^+}^2}{c_3 \cdot c_4} \right) \quad (43)$$

583 where E^θ is the standard reduction potential and has a
584 theoretical value of 1.256 V according to the standard
585 potentials of each part. On the one hand, c_4 , c_5 and c_{H^+}
586 are the molar concentration of VO^{2+} (vanadium(IV)),
587 VO_2^+ (vanadium(V)) and hydrogen ions in the positive
588 half-cell. On the other hand, c_2 , c_3 are the molar con-
589 centration of V^{2+} and V^{3+} , in the negative side.

590 Nevertheless, in most studies (43) is simplified con-
591 sidering only the vanadium species and neglecting the
592 effect of the protons in the positive side, assuming
593 that have not contribution to the equilibrium potential.
594 However, looking at equation (2) it is necessary to con-
595 sider that 2 H^+ are part of the reaction. It has been
596 reported in the literature that if they are neglected, the
597 E^{OCV} value can differ from the real value by 100 mV [8].

598 Another important effect that most of the studies
599 neglect is the Donnan potential, which is related to an
600 unequal distribution of ions in the membrane between
601 different ionic solutions.

602 Accordingly, considering the importance of the pro-
603 tons and having its own electrochemical model that al-
604 lows to understand its evolution over the time, in this
605 work it has been considered to add the Donnan poten-
606 tial to (43) obtaining the following expression for the

607 OCV of a VRFB:

$$E_D^{OCV} = E^\theta + \frac{RT^c}{F} \ln \left(\frac{c_2^c \cdot c_5^c \cdot (c_{H^+}^c)^3}{c_3^c \cdot c_4^c \cdot c_{H^+}^c} \right). \quad (44)$$

608 Considering a stack of N cells, the OCV can be com-
609 puted as the sum of each individual potential. There-
610 fore, assuming that all cells are equal, it can be com-
611 puted by multiplying (44) by N .

612 To analyze the influence of the protons for the com-
613 putation of the OCV, a SOC-OCV profile has been ob-
614 tained through simulation. The SOC has been varied
615 from 5% to 95%, considering a total vanadium concen-
616 tration of $1650 \text{ mol}\cdot\text{m}^{-3}$, and initial proton concentra-
617 tion of 1588 and $2412 \text{ mol}\cdot\text{m}^{-3}$ for the negative and pos-
618 itive sides, respectively. These initial conditions have
619 been selected taking into account the process to obtain
620 the electrolytes used in the experimental setup.

To simplify this first analysis as much as possible,
the evolution of the vanadium and protons concentra-
tion has been directly related to the SOC. Thus, they
are calculated as follows:

$$\begin{aligned} c_2^c &= c_5^c = c_v \cdot \text{SOC} \\ c_3^c &= c_4^c = c_v \cdot (1 - \text{SOC}) \\ c_{H^+}^c &= c_{H^+}^c(0) + c_v \cdot \text{SOC}/2 \\ c_{H^+}^c &= c_{H^+}^c(0) + c_v \cdot \text{SOC}/2. \end{aligned}$$

621 Assuming a constant temperature T^c of 298 K , it is
622 possible to compute the SOC-OCV profile of the gen-
623 eral and most used expression (43), the same one with-
624 out considering the protons, denoted as E_s^{OCV} and the
625 ones considering the Donnan effect presented in (44).

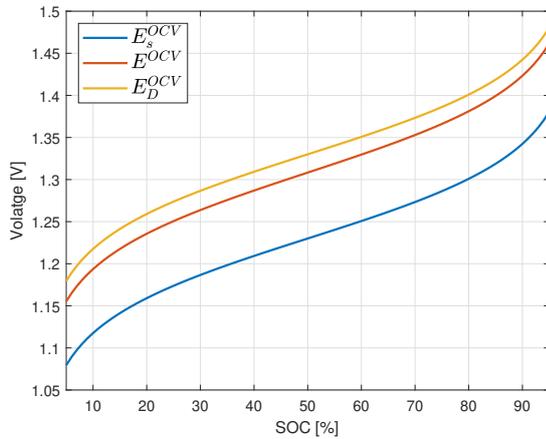


Figure 4: SOC-OCV profiles of a VRFB. E_s^{OCV} : consider-
ing only the vanadium species, E^{OCV} : considering the
 H protons of the polysolyte and E_D^{OCV} : considering the
Donnan effect.

626 Figure 4 shows the three different profiles of the
627 OCV in terms of the SOC, where it is possible to see
628 how there exist a large difference when the protons are
629 not considered, and a small one considering the Don-
630 nan effect. Therefore, it can be concluded that the in-
631 corporation of protons in the VRFB model is a necessary
632 condition for the analysis of the OCV evolution, which
633 is one of the most important variables to be considered.

634 2.5.2 Activation overpotential

635 The activation overpotential, η^a , is defined as the po-
636 tential difference required to overcome the activation
637 energy of the redox reaction to produce a specified cur-
638 rent. It can be computed by means of the Butler-Volmer
639 equation, which has the following expression:

$$j_k = j_k^0 \left(e^{\frac{(1 - \alpha_k)NF\eta_k^a}{RT^c}} - e^{-\frac{\alpha_k NF\eta_k^a}{RT^c}} \right) \quad (45)$$

640 where k is used to distinguish between the positive
641 and negative part, j^0 is the exchange current density
642 in $\text{A}\cdot\text{m}^{-2}$ and α is the charge transfer coefficient. It is
643 important to notice, that the current density has differ-
644 ent sign in the negative and positive parts, so $j_+ = j$ on
645 the positive side, while $j_- = -j$ on the negative one.

646 Therefore, it is possible to calculate the value of the
647 activation overpotential of a cell as:

$$\eta^a = \eta_+^a - \eta_-^a. \quad (46)$$

648 The main problem of the Butler-Volmer expression (45)
649 is that there is not an analytical solution to obtain η^a
650 in terms of the current density. For this reason, most
651 of the studies use other tools to compute its value as a
652 lookup table (LUT) or a function approximation. In this
653 work, we present a hyperbolic sine approximation that
654 is formulated as:

$$\eta^a = \begin{cases} \frac{RT^c}{\alpha F} \cdot \sinh^{-1} \left(\frac{r_j}{2} \right) & r_j \leq \underline{b}_r \\ 0 & \underline{b}_r < r_j < \overline{b}_r \\ \frac{RT^c}{(1 - \alpha)F} \cdot \sinh^{-1} \left(\frac{r_j}{2} \right) & \overline{b}_r \leq r_j \end{cases} \quad (47)$$

655 where r_j is the ratio j/j^0 and \underline{b}_r and \overline{b}_r are the upper
656 and lower bounds, respectively.

657 2.5.3 Ohmic overpotential

658 The ohmic overpotential is directly related to the cell
659 resistance, due to the different materials used, and with
660 the current density j .

661 The cell resistance can be computed as the sum
662 of that imposed by three elements: current collectors,

663 membranes and electrodes. However, in some studies
 664 it has been shown that its value can vary over time and
 665 that depends on whether the system is on a charging
 666 or discharging process. In this work, its value has been
 667 considered constant differentiating between a charging
 668 or a discharging process, according to some results ob-
 669 served in different studies [12].

670 Therefore, the ohmic overpotential η^o can be com-
 671 puted as:

$$\eta^o = \begin{cases} r_d \cdot j & j < 0 \\ r_c \cdot j & j > 0 \end{cases} \quad (48)$$

672 where r_d and r_c represent the cell resistance expressed
 673 in $\Omega \cdot \text{m}^2$.

674 2.5.4 Concentration overpotential

675 The concentration overpotential η^c is related to a poten-
 676 tial that appears inside each electrode, due to the fact
 677 that there exist a difference between the concentration
 678 of species in the bulk electrolyte, compared to the one
 679 on the surface. This difference of concentrations in the
 680 same electrode occurs if the redox reaction of the cell is
 681 much faster than the mass transport. In this scenario,
 682 the redox reaction does not occur ideally, considering
 683 all possible species. On the one hand, it is possible that
 684 some of the reactant species do not reach the reaction,
 685 resulting in a depletion of these species. On the other
 686 hand, considering the product species, a low mass can
 687 cause that these molecules cannot be released from the
 688 reaction, resulting in an accumulation in the surface.

689 It is important to notice, that this phenomenon
 690 must be considered if the model presents mass trans-
 691 port phenomena, as diffusion, migration or convection,
 692 which actually occur in VRFB's. Accordingly, most
 693 works that consider the voltage, do not include the con-
 694 centration overpotential η^c , assuming that there is not
 695 mass transfer effect and therefore, concentrations on
 696 the electrode surface do not differ from the bulk ones.

697 Similarly to the formulation of the activation over-
 698 potential, it is possible to calculate η^c as:

$$\eta^c = \eta_+^c - \eta_-^c \quad (49)$$

699 where η_-^c and η_+^c are the concentration overpotentials
 700 in the negative and positive half-cells, being calculated
 701 by the equations (50) and (51), respectively.

$$\eta_-^c = \frac{RT^c}{F} \ln \left(\frac{1 - \frac{\Delta c}{c_3^c}}{1 + \frac{\Delta c}{c_2^c}} \right) \quad (50)$$

702

$$\eta_+^c = \frac{RT^c}{F} \ln \left(\frac{1 + \frac{\Delta c}{c_5^c}}{1 - \frac{\Delta c}{c_4^c}} \right). \quad (51)$$

703 As can be seen, all overpotentials depend on the cur-
 704 rent, as well as on the species concentration. In order
 705 to analyze their behaviour, they have been computed
 706 numerically considering 3 different charging currents.
 707 Figure 5 presents the profiles of each one of the over-
 708 potentials in terms of SOC and current. It is possible
 709 to see how in the areas close to the maximum charge
 710 and discharge, both activation and concentration over-
 711 potentials have a large contribution in the cell voltage.
 712 Finally, Figure 6 presents the voltage profiles for the
 713 case of a charging process with a constant current of
 714 0.9 A, comparing the OCV with and without the con-
 715 sideration of the Donnan effect and the overpotentials.

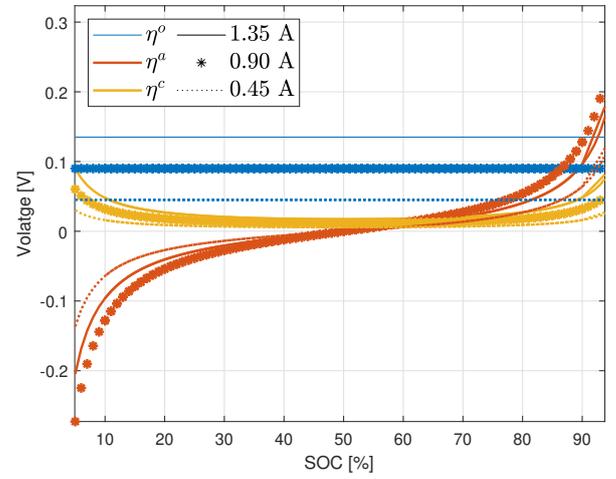


Figure 5: Overpotentials profiles of a VRFB for different values of charging current. η^o : ohmic overpotential, η^a : activation overpotential and η^c : concentration overpotential.

716 3 VRFB CONSERVATION LAWS

717 One of the main characteristics of the electrochemical
 718 model presented here, is that has been formulated ac-
 719 cording to the mass and charge conservation principles.

720 Under these conditions, assuming that the system is
 721 perfectly isolated and has no losses, the total amount of
 722 species will remain constant. The analysis of the mass
 723 and charge conservation can be performed using the
 724 state-space formulation.

725 3.1 Mass conservation

726 Conservation of species in terms of moles for any chem-
 727 ical reaction derives from the mass conservation law.
 728 Thus, it is possible to express the total number of moles
 729 of a specific species as:

$$m_i^k = c_i^k \cdot v^k, \quad (52)$$

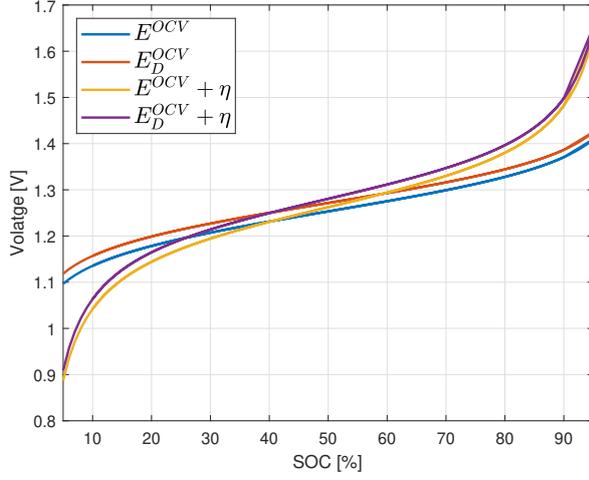


Figure 6: SOC-OCV profiles of a VRFB with and without the consideration of the Donnan effect and the overpotentials.

where m is used to express the mass, defined as the number of moles, and the superscript k represents the place that the species i occupies.

With this notation, it is possible to compute the total number of moles of any RFB, by the following expression:

$$m_t = \sum_{k=c,t} \sum_{i=2}^5 c_i^k \cdot v^k, \quad (53)$$

Therefore, for the VRFB, the mass dynamics for the vanadium electrochemical model can be defined as:

$$\dot{m}_t = v \cdot \dot{x}, \quad (54)$$

where \dot{m}_t defines the dynamics of the number of moles along the time, thus having units of $\text{mol} \cdot \text{s}^{-1}$ and v is the volume vector of a RFB defined as $v = [v^c \ v^c \ v^c \ v^c \ v_-^t \ v_-^t \ v_+^t \ v_+^t]$.

Finally, substituting (15) in the previous expression, it is obtained:

$$\dot{m}_t = v \cdot (\mathbf{A}x + \mathbf{B}_-x \cdot q_- + \mathbf{B}_+x \cdot q_+ + \mathbf{b} \cdot j) = 0, \quad (55)$$

which indicates that the mass conservation law is fulfilled for the vanadium electrochemical model.

Although (55) shows that the total mass is preserved in the global system it does not give information about the behaviour in each half-cell. In this work, the diffusion mechanism has been chosen as an ion crossover effect. Therefore, it is possible to see how the mass changes in each part of the battery due to this effect.

For the case of the negative side, the mass conservation study can be performed with the negolyte volume vector v_- , defined as $v_- = [v^c \ v^c \ 0 \ 0 \ v_-^t \ v_-^t \ 0 \ 0]$.

In this way, it is possible to obtain the following expression:

$$\dot{m}_- = v_- \cdot \mathbf{A}x = \frac{s^m}{w^m} \cdot (-k_2c_2^c - k_3c_3^c + k_4c_4^c + k_5c_5^c). \quad (56)$$

On the other hand, considering the posolyte side, computing the mass differential equation with the volume vector $v_+ = [0 \ 0 \ v^c \ v^c \ 0 \ 0 \ v_+^t \ v_+^t]$, the same results can be extrapolated with opposite sign:

$$\dot{m}_+ = v_+ \cdot \mathbf{A}x = \frac{s^m}{w^m} \cdot (k_2c_2^c + k_3c_3^c - k_4c_4^c - k_5c_5^c). \quad (57)$$

From the expressions presented in (56) and (57) it is possible to see how the sum of both makes the total mass equal to zero, so the mass is conserved in the entire system. However, it can be seen how the mass will be different in each part if the species concentration and the diffusion coefficients are not equal.

The initial species concentration can be known when the electrolytes are prepared. However the exact value of the diffusion coefficients is difficult to know, even more if the dependence on T^c , as expressed in (13), is considered. Due to the impossibility of being able to measure their values, there is no consensus on their values or ratios between them, although the following relationship is satisfied

$$k_2 > k_4 > k_5 > k_3. \quad (58)$$

Using this relationship, it is possible to analyze where and how fast the species are moving through the membrane. In order to compute it, a real Nafion membrane with a thickness of $125 \mu\text{m}$, has been considered. Table 2 summarizes the different diffusion coefficients. They have been computed using the Arrhenius equation (13), considering the pre-factors of the commercial membranes and a constant electrolyte temperature of 298 K. As it can be noticed, the diffusion coefficients follow the relationship presented in (58).

Table 2: Pre-factors and diffusion coefficients of the different vanadium species for a Nafion membrane [27].

Vanadium species	Pre-factor $A \ (\text{m}^2 \cdot \text{s}^{-1})$	Diffusion Coefficient $k \ (\text{m}^2 \cdot \text{s}^{-1})$
V^{2+}	$9.6 \cdot 10^{-9}$	$8.83 \cdot 10^{-12}$
V^{3+}	$3.5 \cdot 10^{-9}$	$3.22 \cdot 10^{-12}$
V^{4+}	$1.1 \cdot 10^{-8}$	$6.83 \cdot 10^{-12}$
V^{5+}	$6.4 \cdot 10^{-9}$	$5.83 \cdot 10^{-12}$

In order to analyze the diffusion mechanism, as well as checking the principle of mass conservation, some simulations have been performed. The considered RFB facility is formed by two tanks of 100 ml of capacity, and a cell with the dimensions shown in Table 3, considering a total concentration of 0.4 M.

Table 3: Model parameters for the conservation of mass study.

Parameter	Value
v_-^t	$1 \cdot 10^{-4} \text{ m}^3$
v_+^t	$1 \cdot 10^{-4} \text{ m}^3$
s^e	$9 \cdot 10^{-4} \text{ m}^2$
s^m	$9 \cdot 10^{-4} \text{ m}^2$
w^e	$5 \cdot 10^{-3} \text{ m}$

791 Considering these characteristics, the total mass of
 792 the system according to expression (53) is 0.004 mol.
 793 Assuming that the system is initially balanced, with
 794 same concentration of $1000 \text{ mol} \cdot \text{m}^{-3}$ for each vanadium
 795 species, a simulation of the diffusion phenomenon has
 796 been performed. Figure 7 shows the evolution of the
 797 four vanadium species, until the equilibrium is reached.
 798 However, it is important to remark that the transport
 799 of species due to diffusion will stop when one of them
 800 disappears, as it happens with c_2 . The evolution of the
 801 system mass is shown in Figure 8 where it can be ob-
 802 served that the total mass m_t is constant according to
 803 the mass conservation principle with an approximate
 value of 0.004 moles.

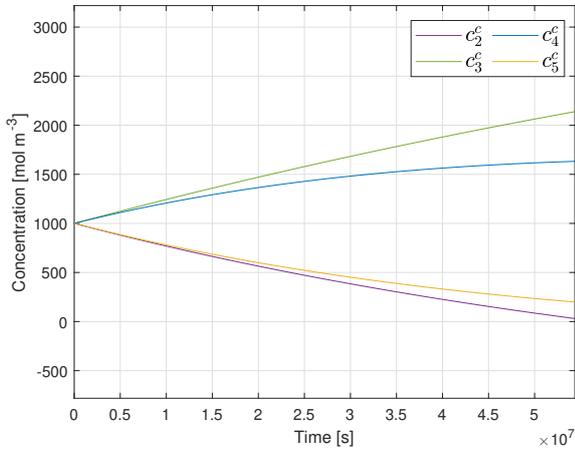


Figure 7: Evolution of vanadium species due to diffusion phenomenon.

804
 805 The negolyte mass, denoted as m_- , increases with
 806 time while the posolyte one, m_+ , decreases, clarifying
 807 the mass transport.

808 3.2 Charge conservation

809 Using the same procedure developed to analyze mass
 810 conservation, it is possible to evaluate if charge con-
 811 servation is also fulfilled. In this case, it is important to
 812 consider that each vanadium species has its own charge
 813 which is directly related to its oxidation state. Thus, the

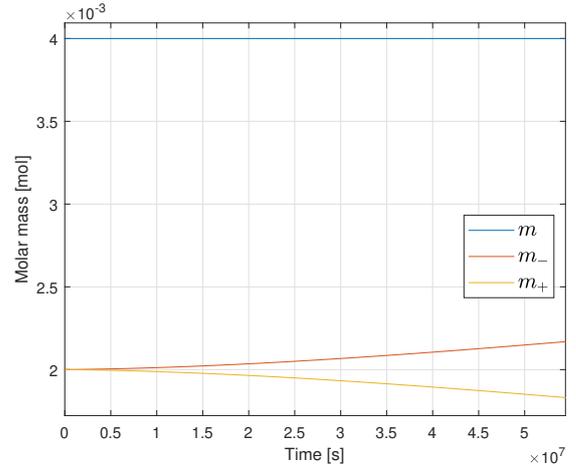


Figure 8: Evolution of total mass m , and posolyte, m_+ , and negolyte m_- due to diffusion phenomenon.

814 expression for the total charge of the system depending
 815 on the number of moles of each species is:

$$c_t = (2c_2^c + 3c_3^c + 2c_4^c + c_5^c)v^c + (2c_2^t + 3c_3^t)v_-^t + (2c_4^t + c_5^t)v_+^t, \quad (59)$$

816 where c_t is used to express the total charge of the sys-
 817 tem, which depends on the valence i of each vanadium
 818 species V^{i+} . Similarly to the analysis performed for the
 819 total number of moles, m_t , it is possible to analyze the
 820 total charge dynamics as:

$$\dot{c}_t = \mathbf{v}_c \cdot \dot{\mathbf{x}}, \quad (60)$$

821 where the volume vector is $\mathbf{v}_c = [2v_-^c, 3v_-^c, 2v_-^c, v_-^c, 2v_+^c, 3v_+^c, 2v_+^c, v_+^c]$, obtaining that its value is zero, which
 822 means that the total charge remains constant inside the
 823 system:

$$\dot{c}_t = \mathbf{v}_c \cdot (\mathbf{A}\mathbf{x} + \mathbf{B}_- \cdot \mathbf{x} \cdot q_- + \mathbf{B}_+ \cdot \mathbf{x} \cdot q_+ + \mathbf{b} \cdot \mathbf{j}) = 0. \quad (61)$$

825 4 EXPERIMENTAL VALIDATION

826 In order to validate the model, some different exper-
 827 iments have been carried out with a real VRFB cell.
 828 A discharge process to calibrate the model parameters
 829 has been considered to later develop a series of charge
 830 and discharge cycles in order to be able to see the effect
 831 of the ion crossover. The change of the mass in each side
 832 of the system, which is translated in the estimation of
 833 the SOH, has been identified.

834 The experimental set-up used in this work consists
 835 of a VRFB single cell developed by the Instituto de Car-
 836 boquímica, two tanks and two peristaltic pumps. The
 837 cell consists of two $7 \times 7 \text{ cm}^2$ porous electrodes separated
 838 by a Nafion-212 membrane. Each tank contains 80 ml

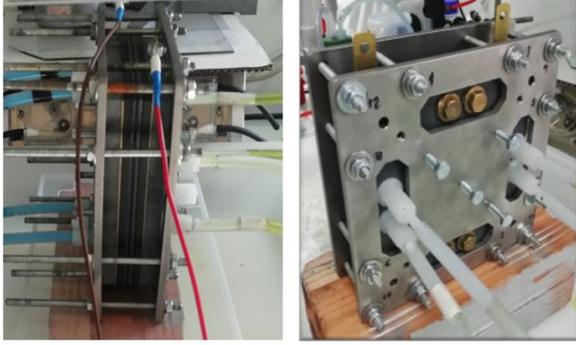


Figure 9: VRFB single cell assembled in the Instituto de Carboquímica center.

839 of 1.6 M vanadium solutions in 2 M sulphuric acid and
840 0.05 M phosphoric acid.

841 The first experiment performed consists of a dis-
842 charge cycle at a constant current of 3 A where the
843 voltage has been measured every 1 second, without a
844 certain idea of the initial conditions in terms of initial
845 species concentration.

846 Despite the fact that many of the variables and pa-
847 rameters of the system are known or can be measured,
848 there is no information available about some others and
849 they need to be estimated. For this purpose, the use
850 of estimation methods is required, and an offline esti-
851 mator based on the PSO technique is proposed. The
852 parameters and variables that have been considered
853 in this work are the initial species concentrations, and
854 the parameters related to the different overpotentials,
855 which are the cell resistance r and the charge transfer
856 coefficients α_+ and α_- . For the case of the standard
857 electrode potential E^θ , it is estimated considering the
858 effect of the Donnan potential on its computation, be-
859 ing denoted as E^{θ^*} .

860 Under the assumption of high flow rates and the
861 mass and charge conservation principles, it is possible
862 to reduce the original system to a 1st order model, being
863 only necessary to calibrate one of the initial concentra-
864 tions, which in this case is the vanadium species c_2 .

865 Using a PSO technique [2], it is possible to estimate
866 the set of parameters $\mathbf{p} = [\alpha_+, \alpha_-, r_d, E^{\theta^*}, c_2(0)]$, that
867 satisfy the following problem:

$$\begin{aligned} \min_{\mathbf{p}} \quad & \sum_{k=1}^N |E(k \cdot T_s) - \hat{E}(k \cdot T_s)| \\ \text{subject to} \quad & \hat{E}(k \cdot T_s) = f(\mathbf{p}) \\ & \mathbf{c}(\mathbf{p}) \leq \mathbf{0} \end{aligned}$$

868 N being the total number of measures that have been
869 taken, considering an equal sample period T_s , k denotes
870 each sample and \mathbf{c} is the constraint set of the unknown

871 parameters. For the experiment performed, the value
872 of T_s is one second and the constraint set is $\mathbf{c}(\mathbf{p}) = [0..1,$
873 $0..1, 0..0.5, 0..2, 0..1600]$.

Table 4: Measured and calibrated VRFB parameters.

Nomenclature	Value	Obtaining
c_t	0.896 mol e ⁻	Measured
m_t	0.264 mol	Measured
m_+	0.132 mol	Measured
T	298 K	Measured
v^c	$7.5 \cdot 10^{-6} \text{ m}^3$	Measured
v_-^t	$8 \cdot 10^{-5} \text{ m}^3$	Measured
v_+^t	$8 \cdot 10^{-5} \text{ m}^3$	Measured
α_+	0.51	Estimated
α_-	0.495	Estimated
$c_2(0)$	$1457 \text{ mol} \cdot \text{m}^{-3}$	Estimated
E^{θ^*}	1.235 V	Estimated
r_d	0.11 Ω	Estimated

873

874 Figure 10 shows the profile of the actual voltage E mea-
875 sured during the discharging profile at constant cur-
876 rent, and the estimated ones with the calibrated param-
877 eters depicted in Table 4. As can be noticed, the agree-
878 ment is reasonably good, presenting a similar voltage
879 performance along all the discharging profile.

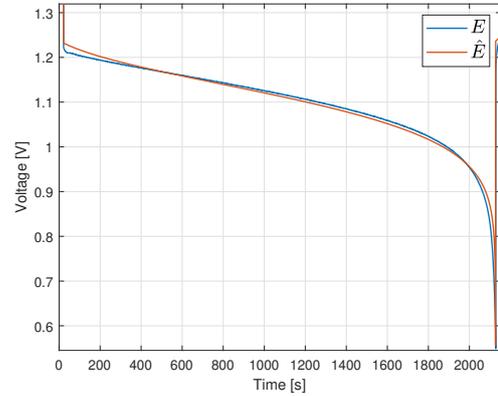


Figure 10: Measured voltage and the estimated ones with the model calibrated.

880 In order to determine if the model presented is
881 able to be used in long-term scenarios when the ion
882 crossover effect is present inside the battery, a series of
883 charge and discharge processes have been carried out.
884 Using the calibrated model and assuming that there is
885 not ion crossover phenomena and that the mass in both
886 sides remains constant and equal to 0.132 mol, the pro-
887 file obtained is the one shown in Figure 11.

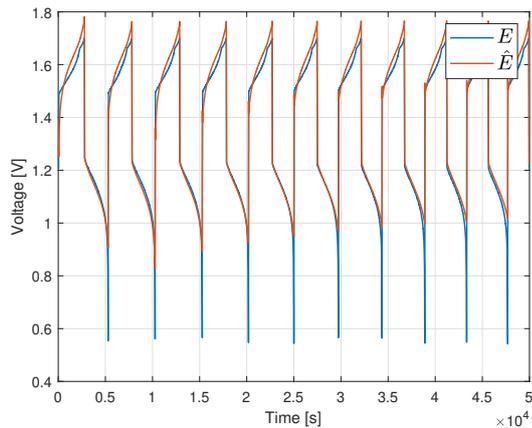


Figure 11: Measured voltage and the estimated ones assuming not change of mass in each side.

888 As can be noticed, there exists a difference between
 889 the measured voltage and the one obtained from the
 890 model. It is also possible to see the effect of the ion
 891 crossover. This experiment consists in charging the sys-
 892 tem with a constant current until it reaches a maximum
 893 voltage value, and discharging it until a minimum one.
 894 As can be observed, initially the duration of the cycles
 895 is approximately 5000 seconds, while at the end of the
 896 profile, it is reduced to around 4300 seconds.

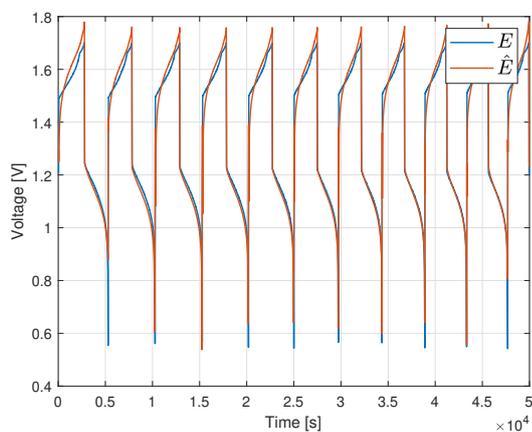


Figure 12: Measured voltage and the estimated ones with a variable m_+ .

897 The difficulty of estimating low vanadium concen-
 898 trations in addition to the existence of side reactions
 899 influences especially in the high voltage region during
 900 charging. This voltage model, based on the Nernst
 901 equation, is highly sensitive to small variations in
 902 vanadium concentration when the SOC is very high or
 903 very low, so it is necessary to improve the concentra-
 904 tion estimation.
 905

906 In order to minimize the error in vanadium concen-
 907 tration, in the implementation of the model with this
 908 experiment, it has been considered that the mass in
 909 the posolyte changes as a ramp, calibrating the value
 910 of the slope in order to estimate the change of mass,
 911 and therefore, the SOH. Performing this analysis using
 912 the same PSO algorithm developed previously and in-
 913 troducing this new parameter for the estimation, the
 914 results obtained are the ones shown in Figure 12.

915 In this case the agreement of the voltage is reason-
 916 ably good compared to the previous ones that did not
 917 consider the mass variation due to crossover. The maxi-
 918 mum relative errors, reached at high voltage during the
 919 charging process, are greater than 50% without con-
 920 sidering mass transport but below 20% in most cycles
 921 when mass transport is considered. The very high de-
 922 viation observed during the short period of time that
 923 the system is above 1.7 V may be related, in addition
 924 to the difficulty of estimating low vanadium concentra-
 925 tions, to the existence of side reactions. Nevertheless,
 926 the mean relative error of the estimation considering
 927 mass transport is around 1.7%. Therefore, and assum-
 928 ing that the system was initially balanced, it is possible
 929 to obtain the profile of the variation of the SOH which
 930 is presented in Figure 13.

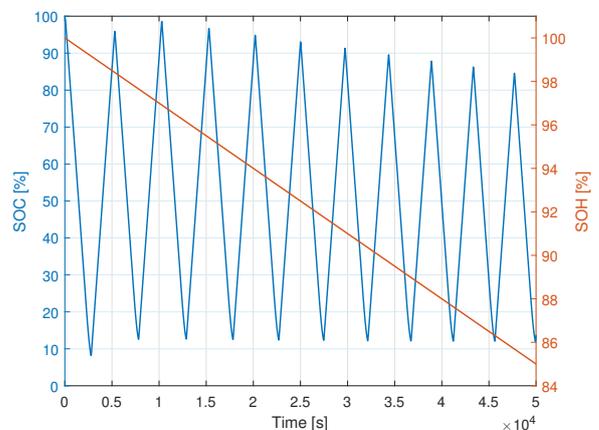


Figure 13: SOC and SOH obtained from the model calibrated considering a variable m_+ .

931 5 Conclusion

932 This work presents a new dynamic model of a VRFB
 933 that allows to consider the most important variables
 934 and effects that occur during the operation of the sys-
 935 tem. The calibration performed with the PSO technique
 936 achieved very accurate results with the data obtained
 937 from a discharge experiment, validating the presented
 938 model that considers different variables. The correct-
 939 ness of the model is guaranteed not only for SOC es-

940 tionation, but it is also able to estimate the SOH when
941 the effect of the ion transport mechanisms is present in
942 a real vanadium redox flow battery. As future improve-
943 ments, it would be necessary to analyze the robustness
944 of the model presented for the case of a stack where the
945 bypass currents can have an important effect.

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